

Variational Hybrid Quantum Algorithms

with Applications to Quantum Chemistry Problems

OsloMet – Autumn School on Quantum Computing

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Chalmers University of Technology

Oslo, November 7, 2023



CHALMERS
UNIVERSITY OF TECHNOLOGY



My Background and Learning Goals

PostDoc at Chalmers University (Oct.2021)
Development of **quantum algorithms** to enable accurate and efficient **quantum chemistry** calculations on current and near-term quantum computers



Gothenburg



140M EUR Research effort for Sweden's Quantum Computing Stack
≈30 PIs, 20 PostDocs and 40 PhDs

NordIQEst:
HPC-QC ecosystem in the Nordics + Estonia



MAX PLANCK INSTITUTE
FOR SOLID STATE RESEARCH

Stuttgart

PhD in **theoretical quantum chemistry** at MPI Stuttgart and University of Stuttgart
Method Development for Quantum Chemistry.
Development of Quantum Monte Carlo methods for strongly correlated electron systems

St. Lorenzen im Paltental



Graz

BSc/MSc Studies in **physics** at TU Graz
Specialization: **Computational/Solid State Physics.**

Take-home messages

- What is quantum chemistry concerned with?
- How can quantum computing help quantum chemistry?
- What are hybrid quantum-classical algorithms?
- How can we simulate chemistry on quantum hardware?
- What is the Variational Quantum Eigensolver (VQE)?

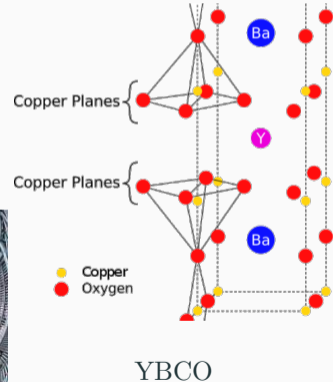
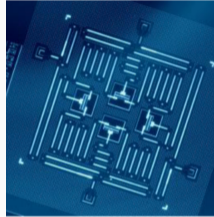
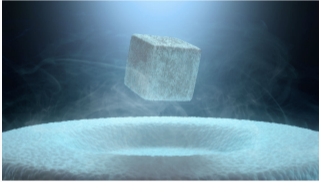
- Motivation – Quantum Chemistry and Electronic Structure Theory
- The Case for Quantum Computing
 - Gate-based Quantum Computing and the Quantum Circuit Model
- How to do quantum chemistry on quantum hardware?
- Quantum Circuit Ansätze and the Variational Quantum Eigensolver
- Applications and Outlook

Motivation – Quantum Chemistry and Electronic Structure Theory

Surprisingly small systems at the center of fascinating
physical and chemical effects

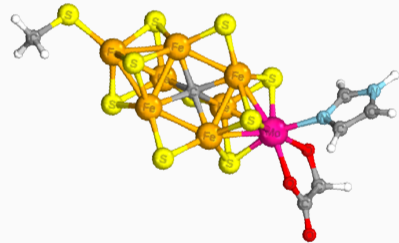
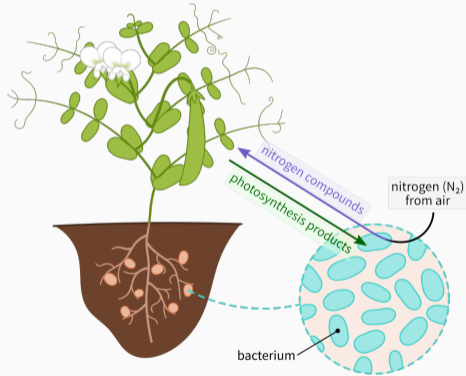
High-temperature superconductivity

Zero electrical resistivity



Nitrogen fixation

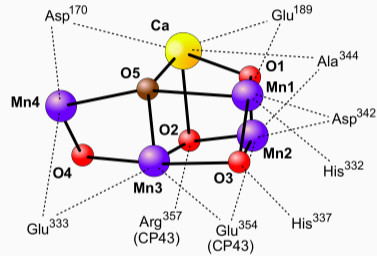
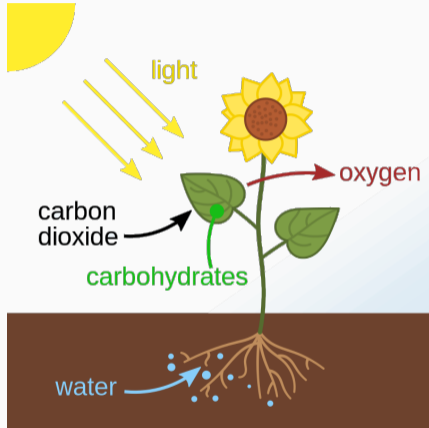
Haber-Bosch process: 1-2% global energy consumption, huge CO₂ emission



Iron-sulfur clusters

⇒ Cheaper and cleaner ammonia production for fertilizers

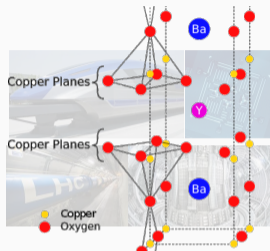
Photosynthesis



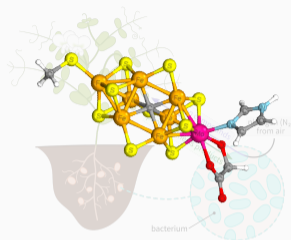
Manganese-Calcium-Oxygen Clusters

Artificial photosynthesis: Carbon capture and hydrogen and oxygen for fuel cell

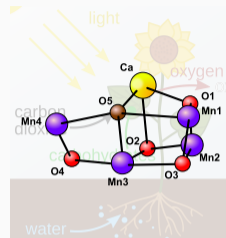
Applications of Quantum Algorithms: *Ab Initio* Quantum Chemistry



YBCO: Unconventional high- T_c superconductivity



FeMoCo: primary cofactor of nitrogenase \rightarrow nitrogen fixation



Manganese-Calcium-Oxygen Clusters: Oxygen evolving clusters in photosystem II

Surprisingly small systems responsible for interesting physical/chemical properties!

Strong electron correlation \Rightarrow challenging systems for computational approaches!

Accurate theoretical understanding at nano-scale for bottom-up materials design!

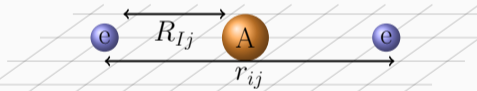
Ab Initio Quantum Chemistry – Electronic Structure Theory

To obtain insight on the **physical** and **chemical properties** (ground- and excited state energies, energy differences, response functions, ...) of quantum systems we need to **solve the Schrödinger equation**:

$$\hat{H} |\Psi\rangle = E |\Psi\rangle,$$

where all necessary information of a quantum system contained in electronic **molecular Hamiltonian**:

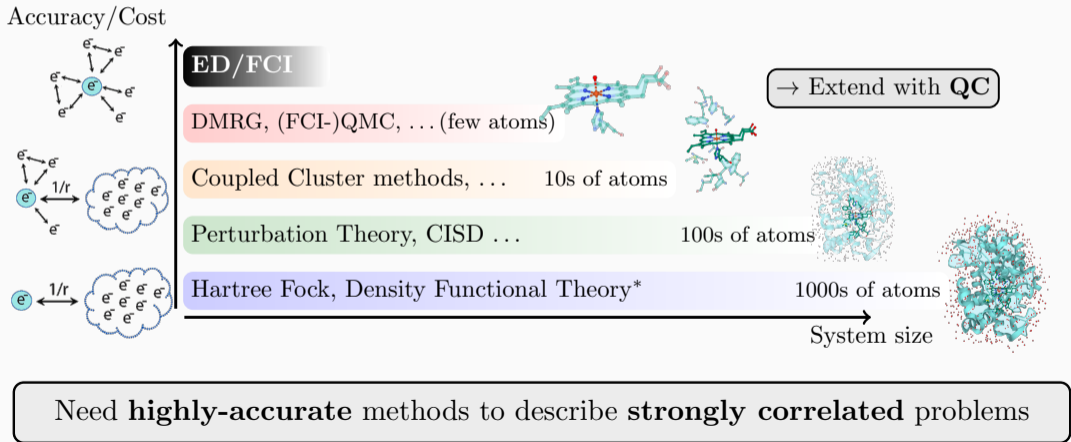
$$\hat{H} = \underbrace{-\frac{1}{2} \sum_i \nabla_{\mathbf{r}_i}^2}_{\text{kinetic energy of } e^-} - \underbrace{\sum_{I,j} \frac{Z_I}{|\mathbf{R}_I - \mathbf{r}_j|}}_{\text{Attr. potential}} + \underbrace{\frac{1}{2} \sum_{i \neq j} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}}_{\text{e}^- - \text{e}^- \text{ repulsion}}$$



Coulomb repulsion correlates all electrons of a system → analytic solution too complex
→ **approximations and computational approaches**

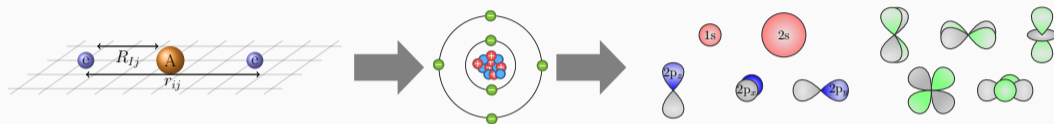
Accuracy and cost – scaling and hierarchy of methods

Depending how accurately we treat correlation: various methods and **levels of theory** to solve the Schrödinger equation



Ab Initio Quantum Chemistry – Electronic Structure Theory

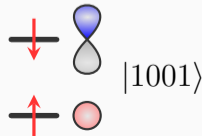
We have to choose a **basis/orbitals** we perform our calculations in! In quantum chemistry: starting orbitals are most often “atomic-like” orbitals (for each atom):



Second quantized form of molecular Hamiltonian:

$$\hat{H} = \underbrace{\sum_{i,j} t_{ij} \sum_{\sigma=\uparrow,\downarrow} a_{i,\sigma}^\dagger a_{j,\sigma}}_{\text{kinetic/hopping term}} + \underbrace{\sum_{i,j,k,l} V_{ijkl} \sum_{\sigma,\tau=\uparrow,\downarrow} a_{i,\sigma}^\dagger a_{j,\tau}^\dagger a_{l,\tau} a_{k,\sigma}}_{e^- - e^- \text{ interaction term}}$$

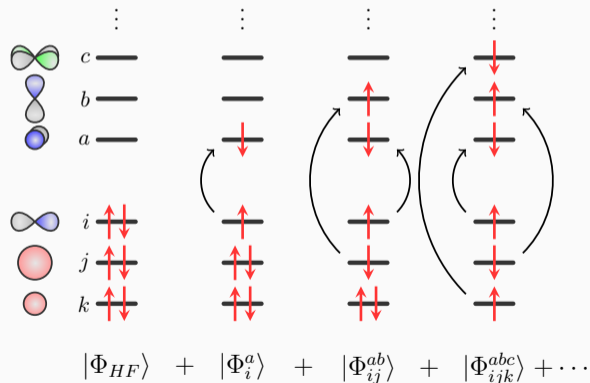
$a_{i,\sigma}^{(\dagger)}$ annihilates(creates) an electron with spin $\sigma = \{\uparrow, \downarrow\}$ in orbital i . e.g. for 2 electrons in 2 orbitals (=4 spin-orbitals): $a_{1,\uparrow}^\dagger a_{2,\downarrow}^\dagger |\text{vac}\rangle = |1001\rangle$:



Exponential scaling of Full Configuration Interaction

FCI \Rightarrow exact solution in a given basis: linear combination of determinants

$$|\Psi\rangle = |\Phi_{HF}\rangle + \sum_i c_i |\Phi_i\rangle$$



All possible excitations from HF determinant

Number of possible states for given number of electrons, N , and orbitals, n , $\sim \binom{N}{n}$

Mol.	#orbitals	#electrons	#states
H ₂	2	2	4
LiH	4	4	36
Be ₂	8	8	4900
H ₂ O	12	12	$\sim 8 \cdot 10^5$
C ₂ H ₄	16	16	$\sim 16 \cdot 10^6$
F₂	18	18	$\sim 2 \cdot 10^9$

The Case for Quantum Computing

Classical bit

0

1

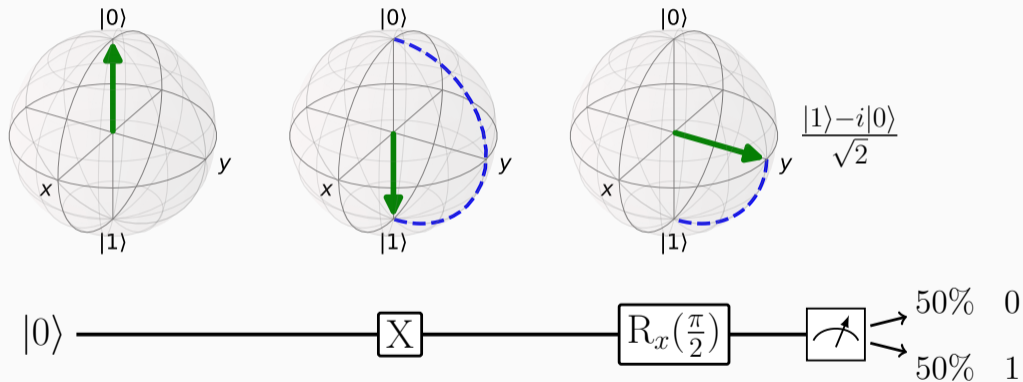
Quantum bit = qubit

$$a |0\rangle + b |1\rangle$$

Quantum bit = qubit

$$a |0\rangle + b |1\rangle$$
$$|a|^2 + |b|^2 = 1$$

Qubits – Bloch Sphere



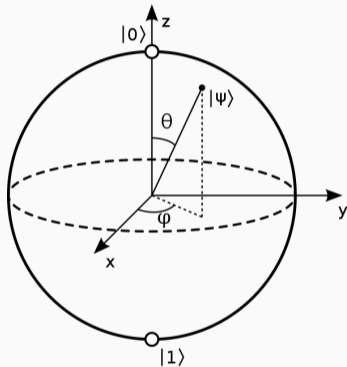
Parametrized gates

For flexibility, we need **parametrized** gates, e.g. rotation around axis:

$$R_X(\theta) = \exp(-i\frac{\theta}{2}\hat{X}) = \begin{pmatrix} \cos(\theta/2) & -i\sin(\theta/2) \\ -i\sin(\theta/2) & \cos(\theta/2) \end{pmatrix}$$

$$R_Y(\theta) = \exp(-i\frac{\theta}{2}\hat{Y}) = \begin{pmatrix} \cos(\theta/2) & -\sin(\theta/2) \\ \sin(\theta/2) & \cos(\theta/2) \end{pmatrix}$$

$$R_Z(\varphi) = \exp(-i\frac{\varphi}{2}\hat{Z}) = \begin{pmatrix} \exp(-i\varphi/2) & 0 \\ 0 & \exp(i\varphi/2) \end{pmatrix}$$



with the Pauli matrices: $\hat{X} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\hat{Y} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\hat{Z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

Multiple Qubits

Bringing **two** qubits together:

$$|\Psi\rangle = \overbrace{(|0\rangle + |1\rangle)}^{\text{qubit 1}} \otimes \overbrace{(|0\rangle + |1\rangle)}^{\text{qubit 2}} = |00\rangle + |01\rangle + |10\rangle + |11\rangle \quad 4 \text{ states}$$

Three qubits:

$$\begin{aligned} |\Psi\rangle &= \overbrace{(|0\rangle + |1\rangle)}^{\text{qubit 1}} \otimes \overbrace{(|0\rangle + |1\rangle)}^{\text{qubit 2}} \otimes \overbrace{(|0\rangle + |1\rangle)}^{\text{qubit 3}} \\ &= |000\rangle + |001\rangle + |010\rangle + |100\rangle + |011\rangle + |101\rangle + |110\rangle + |111\rangle \quad 8 \text{ states} \end{aligned}$$

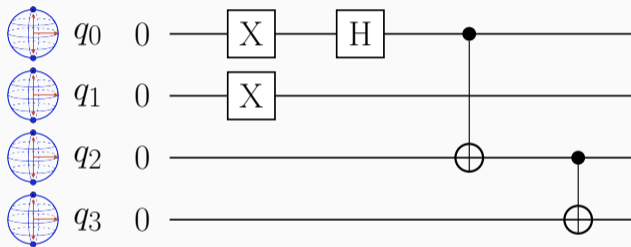
n qubits can encode exponentially many (2^n) states.

→ Need new **quantum algorithms** to use this potential advantage!

≈ 40 qubits enough to encode the $\sim 2 \cdot 10^9$ states of $F_2!$

Quantum circuit model

Circuit model: Lines represent qubits and similar to classical circuits (AND, OR, ...) we can act with **operations/gates** on one (rotations) or **multiple qubits** (CNOT, ...)



Hadamard gate (H):

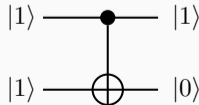
$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

Controlled NOT (CNOT):

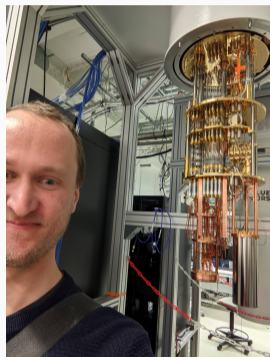
target qubit (\oplus) is inverted if control qubit (\cdot) is in $|1\rangle$ state

Exercise: what is the state $|\Phi\rangle$ at the end of the above circuit?

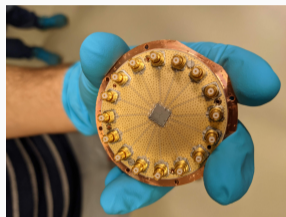
Parametrized gates ($R_x(\theta), \dots$) and multi-qubit gates allow us to prepare flexible **entangled** (non-classical) states, $|\Psi(\theta)\rangle = \hat{U}(\theta)|0\rangle$.



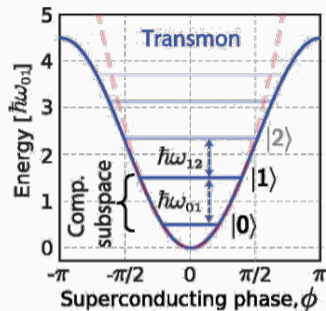
Noisy intermediate-scale quantum - NISQ



@ Chalmers



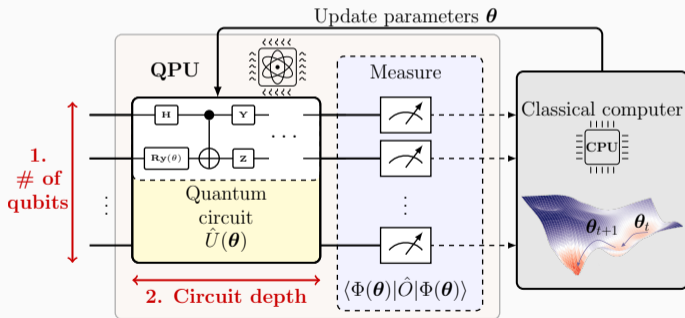
3 qubit device @Chalmers



Qubits rely on quantum effects → very fragile, easily influenced by environmental effects/noise. Need to isolate and cool them close to absolute zero! Only few of them...

Hybrid quantum-classical approach

Current quantum hardware has many flaws: **noise, decoherence and limited number of qubits** → Hybrid quantum-classical approach



Use pros of both CPUs and QPUs:

- Use **short-depth quantum circuits** that fit current hardware
- Can **improve on classical estimates** by non-classical states
- Store quantum state with **exponentially fewer resources**
- Use CPU to i.e. optimize gate parameters, θ

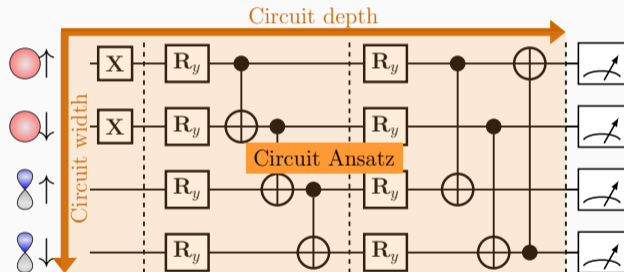
How to do quantum chemistry on
quantum hardware?

Quantum Chemistry on Quantum Computers

1. Map fermionic Hamiltonian/basis functions onto quantum hardware/qubits*
2. Choose an 'Ansatz'
3. Use quantum algorithms for ground-, excited states, dynamics, ...

a) Prepare an initial state

$$|\Phi_0\rangle = \begin{pmatrix} 1 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$



b) Perform **unitary** operations of the quantum algorithm on the qubits

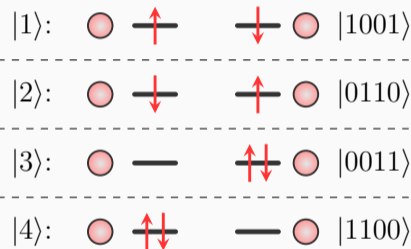
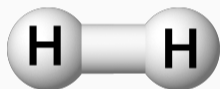
c) Measure observables of interest, ($\langle \hat{H} \rangle$)

$$|\Phi\rangle = \hat{U} |\Phi_0\rangle = a_1 \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} + a_2 \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} + \dots + a_{2^N} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix}$$

* Jordan-Wigner, Bravyi-Kitaev, (Ann. Phys. **298**, 210 (2002)), Parity encoding ...

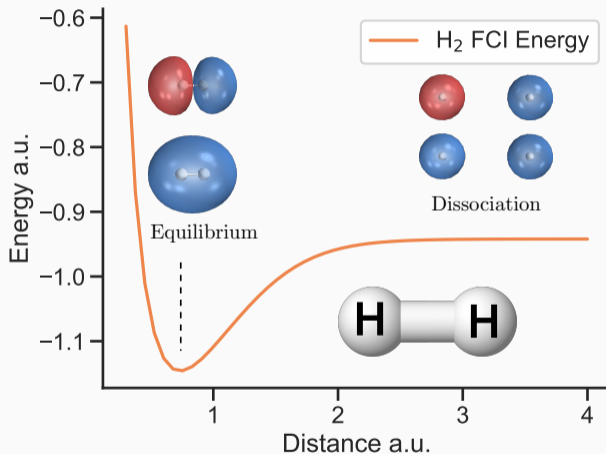
Example: Classic solution of hydrogen molecule – H_2

Hydrogen molecule in a “minimal” basis set: 1s orbital for each hydrogen.
4-dimensional Hilbert space for two electrons in two orbitals (4 spin-orbitals).



Solution: Construct matrix representation of quantum chemistry Hamiltonian in this basis, $H_{ij} = \langle i | \hat{H} | j \rangle$, \rightarrow diagonalize \rightarrow exact solution in given basis

Example: Hydrogen molecule – H₂ – PES



Groundstate wavefunction at equilibrium: $|\Psi_0\rangle = 0.995 |1100\rangle - 0.105 |0011\rangle$

Groundstate wavefunction at dissociation: $|\Psi_0\rangle = \frac{1}{\sqrt{2}} (|1100\rangle - |0011\rangle)$

Encoding the problem on a quantum computer

Qubits can quite naturally store the occupation of an spin-orbital:

$|0\rangle$ -state/ $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$: empty, $|1\rangle$ -state/ $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$: occupied

Also: action of creation and annihilation operators representable by Pauli matrices:

$$\hat{a}^\dagger |0\rangle = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\hat{X} - i\hat{Y}}{2} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |1\rangle$$
$$\hat{a} |1\rangle = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{\hat{X} + i\hat{Y}}{2} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |0\rangle,$$

with the Pauli matrices

$$\hat{X} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \hat{Y} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \hat{Z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

(Anti-)Symmetric wave functions – Pauli exclusion principle

However:

- For Fermions, the Pauli exclusion principle requires the wavefunction to be **anti-symmetric** under the exchange of two particles:

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_i, \dots, \mathbf{r}_j, \dots, \mathbf{r}_N) = -\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_j, \dots, \mathbf{r}_i, \dots, \mathbf{r}_N)$$

- While Bosonic wave functions are **symmetric** under the exchange of particles

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_i, \dots, \mathbf{r}_j, \dots, \mathbf{r}_N) = +\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_j, \dots, \mathbf{r}_i, \dots, \mathbf{r}_N)$$

Anti-symmetry/Anti-commutation relations

Electrons are **indistinguishable Fermions**, with anti-symmetric wavefunction and anti-commuting creation/annihilation operators:

$$\{a_i^\dagger, a_j^\dagger\} = 0, \quad a_i^\dagger a_j^\dagger = -a_j^\dagger a_i^\dagger$$

Unlike **individually addressable** qubits, with Pauli operators:

$$\left(\frac{X_i - iY_i}{2}\right) \left(\frac{X_j - iY_j}{2}\right) = + \left(\frac{X_j - iY_j}{2}\right) \left(\frac{X_i - iY_i}{2}\right)$$

Task: We need to map the **fermionic** Hamiltonian to a **qubit** Hamiltonian in terms of Pauli operators:

$$\hat{H}_f = \sum V_{ijkl} a_{i,\sigma}^\dagger a_{j,\tau}^\dagger a_{l,\tau} a_{k,\sigma} \quad \Rightarrow \quad \hat{H}_q = \sum_i c_i \hat{P}_i$$

Jordan Wigner Mapping

Note that the following Pauli operators anti-commute:

$$\hat{Z}\hat{X} = -\hat{X}\hat{Z}, \quad \hat{Z}\hat{Y} = -\hat{Y}\hat{Z}$$

Exercise: convince yourself of equation above

Jordan-Wigner encoding:

$$a_1^\dagger = \frac{X_1 - iY_1}{2}$$

$$a_2^\dagger = Z_1 \left(\frac{X_2 - iY_2}{2} \right)$$

$$a_3^\dagger = Z_1 Z_2 \left(\frac{X_3 - iY_3}{2} \right)$$

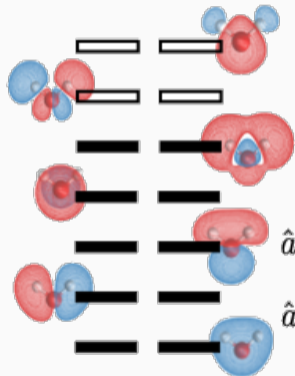
$$a_i^\dagger = \bigotimes_{j=1}^{i-1} Z_j \left(\frac{X_i - iY_i}{2} \right)$$

- Fix anti-symmetry by tracking parity/phase before each creation/annihilation operator with \hat{Z}_i
- **Exercise:** Demonstrate anti-symmetry of JW-encoded $a_3^\dagger a_2^\dagger = -a_2^\dagger a_3^\dagger!$

JW Mapping – Recap

Fermions

$$\{a_i, a_i\} = 0, \{a_i^\dagger, a_i^\dagger\} = 0, \{a_i, a_i^\dagger\} = \delta_{i,j}$$



mapping

Jordan-Wigner

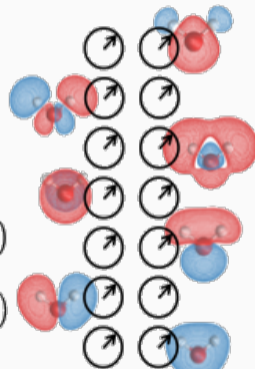
$$\hat{a}_j = \bigotimes_{i=1}^{j-1} \hat{\sigma}_i^z \otimes (\hat{\sigma}_j^x + i\hat{\sigma}_j^y)$$

$$\hat{a}_j^\dagger = \bigotimes_{i=1}^{j-1} \hat{\sigma}_i^z \otimes (\hat{\sigma}_j^x - i\hat{\sigma}_j^y)$$

$$\hat{H}_{elec} = \sum_{pq} h_{pq} \hat{a}_p^\dagger \hat{a}_q + \frac{1}{2} \sum_{pqrs} h_{pqrs} \hat{a}_p^\dagger \hat{a}_q^\dagger \hat{a}_r \hat{a}_s$$

Spins

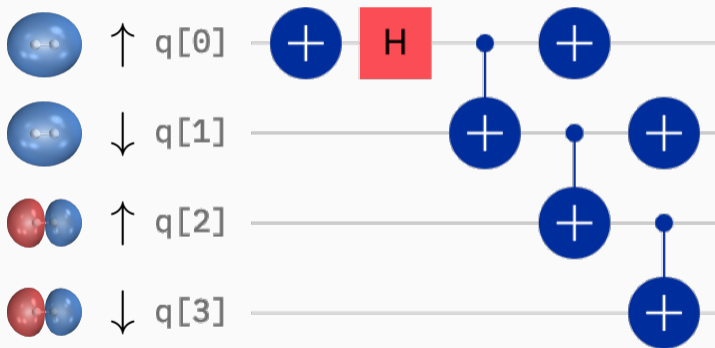
$$[\sigma_i, \sigma_i] = 0, [\sigma_i^\dagger, \sigma_i^\dagger] = 0, [\sigma_i, \sigma_j^\dagger] = \delta_{i,j}$$



$$\hat{H}_{elec} = \sum_i c_i \hat{P}_i$$

Re: H_2 at dissociation

What is the state of this quantum circuit?



Try: [IBM Quantum Composer](#)

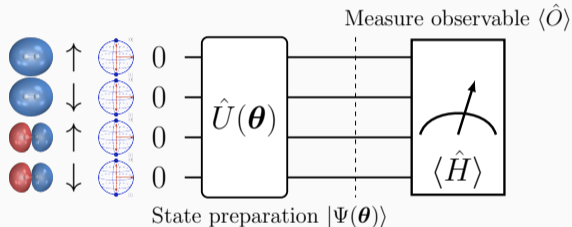
More general quantum circuits – Ansätze

H₂ groundstate wavefunction at equilibrium:

$$|\Psi_0\rangle = 0.995 |1100\rangle - 0.105 |0011\rangle + 0 \cdot |1001\rangle + 0 \cdot |0110\rangle$$

Need a more general “Ansatz”, $\hat{U}(\boldsymbol{\theta}) |\mathbf{0}\rangle$:

$$\begin{aligned} |\Psi(\boldsymbol{\theta})\rangle &= \hat{U}(\boldsymbol{\theta}) |0000\rangle \\ &\stackrel{!}{=} c_1 |1100\rangle + c_2 |0011\rangle \\ &\quad + c_3 |1001\rangle + c_4 |0110\rangle \end{aligned}$$



Given a general quantum Ansatz $\hat{U}(\boldsymbol{\theta})$: We need algorithms to find the optimal gate parameters $\boldsymbol{\theta}$! The variational quantum eigensolver (VQE) can find the **most optimal** parameters/angles $\boldsymbol{\theta}$ to minimize the energy expectation value:

$$E(\boldsymbol{\theta}) = \min_{\boldsymbol{\theta}} \langle \mathbf{0} | \hat{U}^\dagger(\boldsymbol{\theta}) \hat{H} \hat{U}(\boldsymbol{\theta}) | \mathbf{0} \rangle$$

Quantum Circuit Ansätze and the Variational Quantum Eigensolver

Ansatz for the quantum chemistry wavefunction

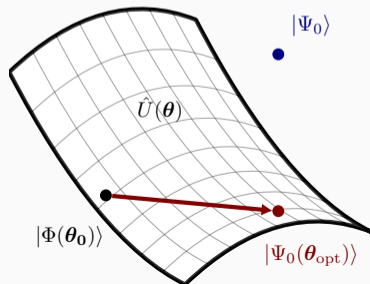
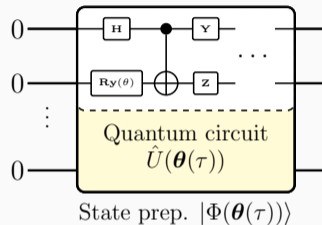
In general: an **Ansatz** is a quantum circuit with parametrized gates

Desired in Ansätze:

- Expressive – spans large and correct portion of Hilbert space
 - Small number of qubits
 - Short-depth
- **Hardware efficient**

Consideration due to noise and coherence times:

- Circuit depth
- Circuit connectivity
- Number of parameters
- Number of 2-qubit gates
- Gate types (Native to hardware?)

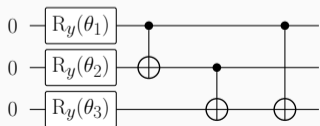


Ansatz for the quantum chemistry wavefunction: $|\Psi(\theta)\rangle = \hat{U}(\theta) |0\rangle$

Hardware efficient Ansätze:

- Gates directly tailored for the specific quantum device
- Spans a very large portion of the Fock space (inefficient)
- Large number of parameters, hard to optimize, “barren plateaus”

$$|\Psi(\theta)\rangle = \prod_i^d \left[\hat{U}_{ent} \hat{U}_{rot}(\{\theta_i\}) \right] |\psi_{init}\rangle$$

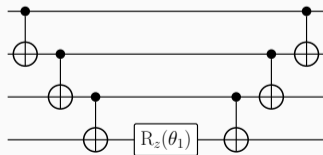


⇒ Adaptive Ansätze*

Chemically/Physically motivated:

- Chemistry-inspired exponential Ansatz: *Unitary coupled cluster Ansatz*
- Encode excitations of electrons between different orbitals, $|1100\rangle \rightarrow |0011\rangle$
- Not hardware efficient, deep circuits
- Less parameters, easier to optimize

$$|\Psi(\theta)\rangle = e^{\hat{T}(\theta) - \hat{T}^\dagger(\theta)} |\psi_{HF}\rangle$$



Ground state, $|\Psi_0\rangle$, is fundamental in quantum chemistry and electronic structure theory \rightarrow used to calculate all sort of properties, like reaction rates and reaction pathways

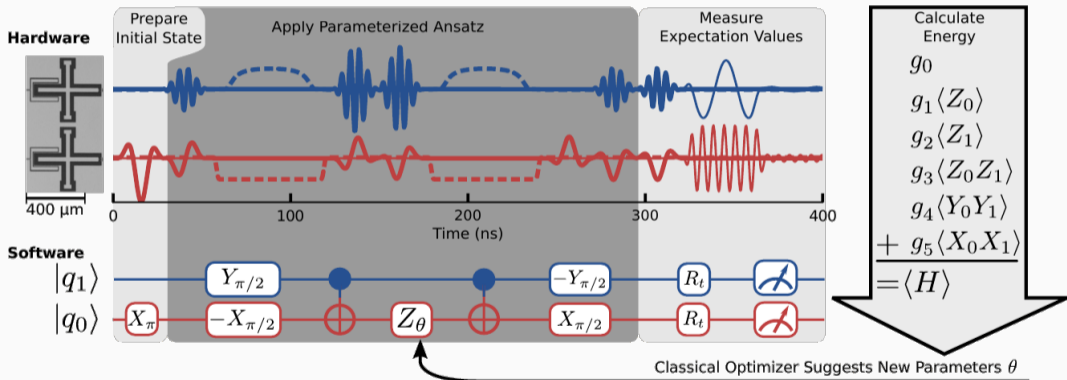
Our goal is to estimate: $\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle \Rightarrow \langle \Psi_0 | \hat{H} | \Psi_0 \rangle = E_0$

Variational principle: an arbitrary state, $|\Psi(\boldsymbol{\theta})\rangle$, the expectation value of \hat{H} , will be an upper bound to E_0

$$\langle \Psi(\boldsymbol{\theta}) | \hat{H} | \Psi(\boldsymbol{\theta}) \rangle = E(\boldsymbol{\theta}) \geq E_0$$

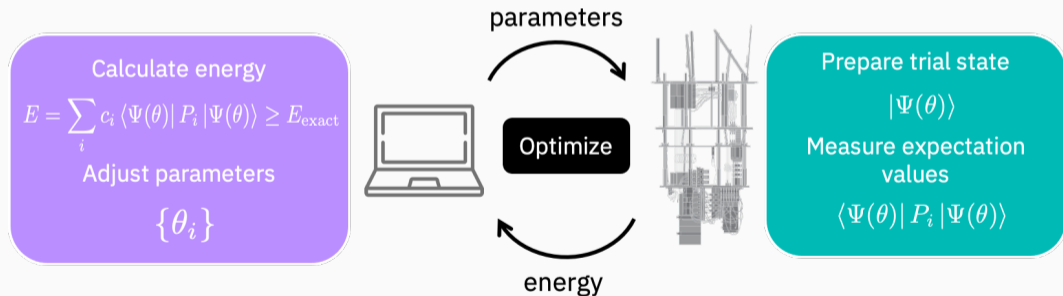
Variational Quantum Eigensolver

Leverage pros of both classical and quantum computers:



Variational Quantum Eigensolver

VQE: Efficiently prepare and encode $|\Psi(\theta)\rangle$ with a suitable *Ansatz*, $\hat{U}(\theta)$, on **quantum hardware** and measure the expectation values of “Pauli strings”, \hat{P}_i . Reconstruct the energy and update the parameters (with some optimizer) on a **classical computer**. Repeat until convergence of $E(\theta) = \langle \Psi(\theta) | \hat{H} | \Psi(\theta) \rangle$



Applications and Outlook

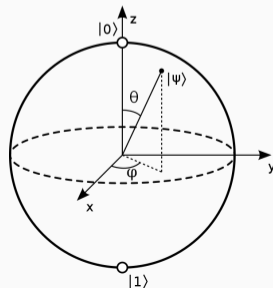
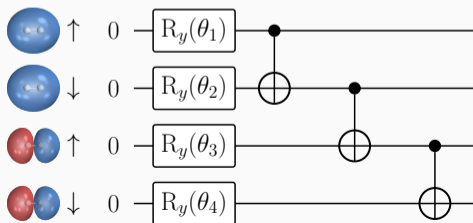
H₂ – Hardware efficient Ansatz

H₂ groundstate wavefunction at equilibrium:

$$|\Psi_0\rangle = 0.995 |1100\rangle - 0.105 |0011\rangle + 0 \cdot |1001\rangle + 0 \cdot |0110\rangle$$

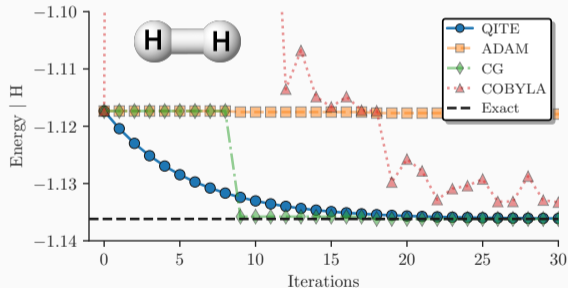
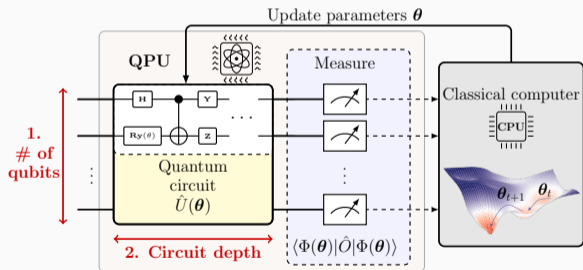
Need a general “Ansatz”, $\hat{U}(\boldsymbol{\theta}) |\mathbf{0}\rangle$:

$$|\Psi(\boldsymbol{\theta})\rangle = \hat{U}(\boldsymbol{\theta}) |0000\rangle \stackrel{!}{=} c_1(\boldsymbol{\theta}) |1100\rangle + c_2(\boldsymbol{\theta}) |0011\rangle + c_3(\boldsymbol{\theta}) |1001\rangle + c_4(\boldsymbol{\theta}) |0110\rangle$$



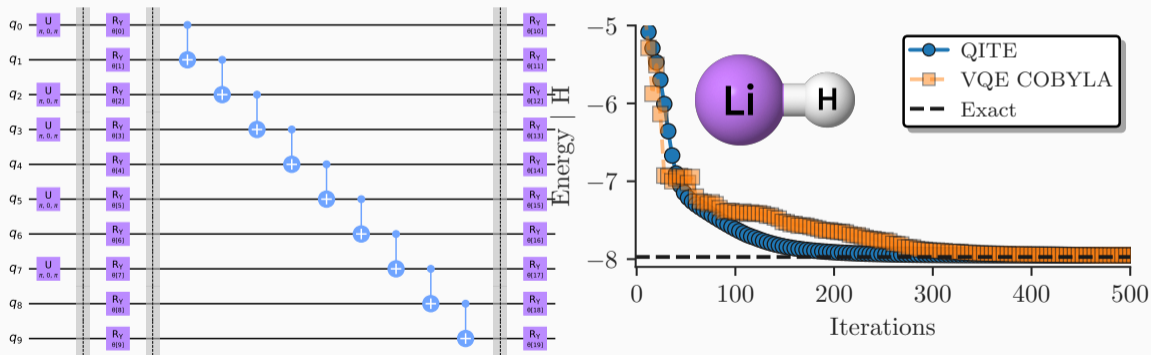
H₂ VQE Statevector simulation

H₂ at 0.7 Å in a STO-6G basis using 4 qubits: Default Qiskit VQE settings for different types of optimizers, no noise!



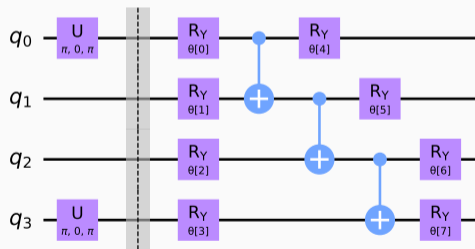
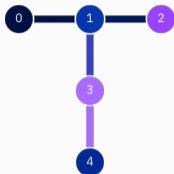
Jupyter Notebook Example

LiH – minimal basis – 1.6\AA – parity mapping – HEA Ansatz

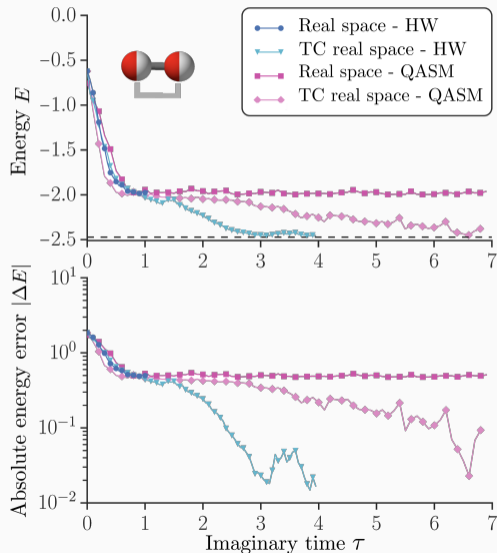


Actual experimental results for the Hubbard model on ibmq_lima

- 2-site Hubbard model
- On ibmq_lima



Hardware-efficient RY Ansatz



Reference-state Error Mitigation

JCTC
Journal of Chemical Theory and Computation

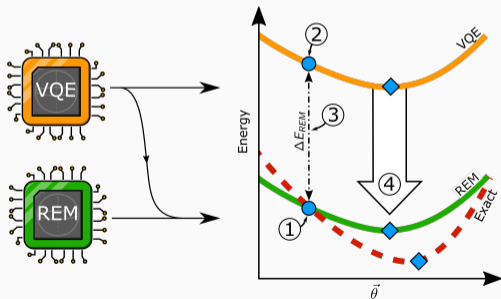
pubs.acs.org/JCTC

Article

Reference-State Error Mitigation: A Strategy for High Accuracy Quantum Computation of Chemistry

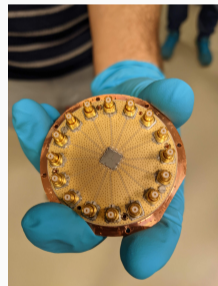
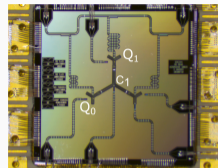
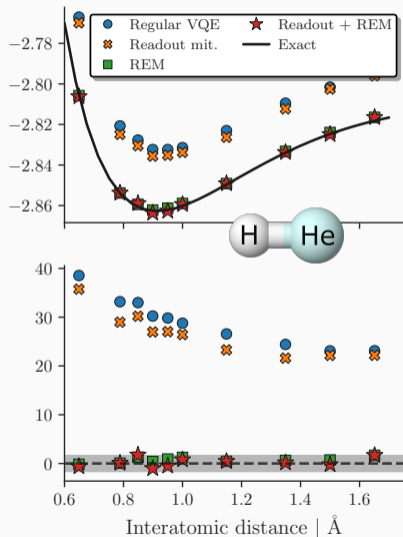
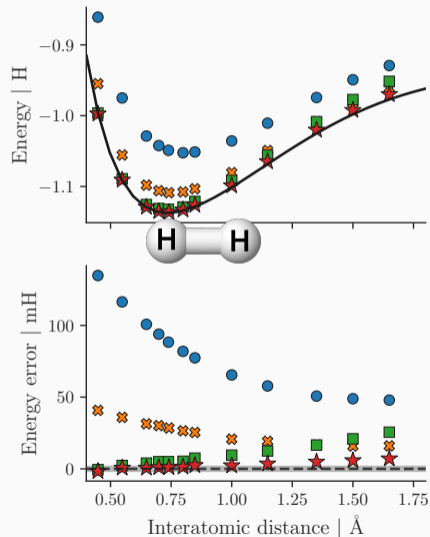
Phalgun Lolur,^{||} Mårten Skogh,^{||} Werner Dobrautz, Christopher Warren, Janka Biznárová, Amr Osman, Giovanna Tancredi, Göran Wendin, Jonas Bylander, and Martin Rahm*

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1. Classically compute exact reference energy (i.e. Hartree-Fock), $E(\theta_{\text{ref}})$
2. Measure reference energy on noisy device $\mathcal{E}(\theta_{\text{ref}})$, with reference parameters θ_{ref}
3. Calculate REM correction:
$$\Delta E_{\text{REM}} = \mathcal{E}(\theta_{\text{ref}}) - E(\theta_{\text{ref}})$$
4. Correct final VQE energy with REM correction

Reference-state Error Mitigation



- Noise!
- Scaling up!
- Calculations on real-hardware over the cloud
- Error mitigation: Reference-State Error Mitigation: A Strategy for High Accuracy Quantum Computation of Chemistry, with calculations/experiments on Chalmers devices
- Resource reduction: Phys. Rev. Research 5, 023174, arXiv:2303.02007
- Many more...

Algorithms on quantum hardware

On near-term intermediate-scale quantum (NISQ) hardware (low number of qubits and short circuits): **hybrid quantum-classical algorithms**

- Quantum Approximate Optimization algorithm (QAOA)
- Quantum (Imaginary) Time Evolution
- **Variational Quantum Eigensolver (VQE)**
- ...

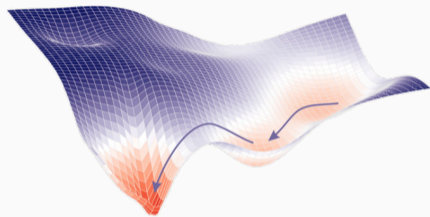
There is a variety of quantum algorithms for different kind of problems, most are for future **fault-tolerant** quantum hardware with many qubits and deep circuits.

- Shor's algorithm – Encryption
- Grover's algorithm – Database search
- Quantum Fourier Transformation
- **Quantum Phase Estimation**

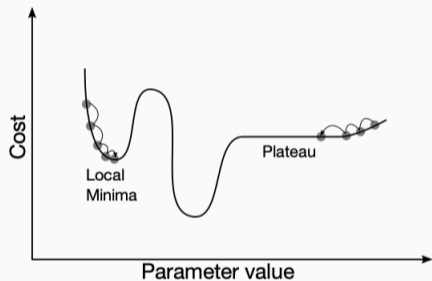
Thank you for your attention!

Variational Quantum Eigensolver – VQE

Main hybrid quantum-classical approach, many advantages (see other talks).



Parameter landscape



Problems, e.g. difficult classical optimizations of parameters θ (**local minima**, **barren plateaus**) and relies on variational principle: $E_0 \leq \min_{\theta} E(\theta) = \langle \Psi(\theta) | \hat{H} | \Psi(\theta) \rangle$
→ not applicable to open quantum system and transport problems (non-Hermitian Hamiltonians)